

Article

Temperature and Conductivity as Indicators of the Morphology and Activity of a Submarine Volcano: Avyssos (Nisyros) in the South Aegean Sea, Greece

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Abstract: The morphology and the activity of a submarine caldera, Avyssos, at the northern part of Nisyros volcano in the South Aegean Sea (Greece), has been studied by means of remotely-operated underwater vehicle dives. The recorded time series of temperature and conductivity over the submarine volcano have been analyzed in terms of the generalized moments method. The findings of the mathematical analysis shed light on the volcanic activity, but also on the morphology (shape) of the submarine volcano. The conductivity time series indicates that the volcano is at rest, in agreement with other types of observations. On the other hand, temperature fluctuations, which in general describe a multifractal process, show that the submarine caldera operates as an open system that interacts with its surroundings. This type of analysis can be used as an indicator for the state of activity and the morphological structure (closed or open system) of a submarine volcano.

Keywords: generalized moments method; time series; submarine volcano; Nisyros

1. Introduction

Nisyros island is an active volcano at the eastern edge of the Aegean within the Hellenic Volcanic arc and trench system along the convergence zone of the Eurasian and African plates [1,2]. Seismic activity during 1995–1998 affected the island with damage along the western edge of the island of Mandraki (located in NW Nisyros), related to the Mandraki fault, which extends northwards under the sea [3]. Several types of studies have been conducted on Nisyros since then, focusing on seismicity [4], tectonics [5], petrology [6], geochemistry [7] and more. Remotely-Operated Underwater Vehicle (ROV)-based studies have been focused on the geomorphology of the underwater features around Nisyros and explored the volcanic field offshore of the island, mainly along the Nisyros-Yali-Strongyli area [8]. Such information is available for the neighboring Avyssos underwater caldera (see Figure 1) at the northeastern part of the volcanic islet of Strongyli, north of Nisyros volcano. The study of underwater volcanic activity of the area can provide important information on the properties of the volcanic system, especially in comparison to the shallow active submarine volcano Kolumbo, NE of Santorini [9–12]. Kolumbo was studied extensively in the summer and fall of 2010 and the summer of 2011 using the ROVs “Hercules” and “Argus” of E/V “Nautilus”. Among others, a dedicated cruise in the 2010 campaign focused on exploring the unknown volcano-tectonic characteristics of the Nisyros-Yali-Kos region and the associated volcanic field around Nisyros. Both Nisyros and Santorini belong to the same volcanic arc; both have exhibited unrest in the recent past [4,13], and

both present potential sources of geohazards [4,8]. However, these underwater volcanoes have quite different geomorphologies: Kolumbo in Santorini is a bowl-shaped volcano with an elevation difference between its rim and bottom of about 400 m (the maximum depth is 504 m) [10], while the Avysson caldera's dimensions are 3 km in the NW-SE direction and 4 km in the NE-SW direction with a maximum depth of 673 m [8]. In addition, there is a hill of about 1 km in length at the center of Avysson with a positive relief of about 60–70 m above the flat sea bottom (Figure 1). These differences may cause different behaviors in the recorded physical and chemical data.

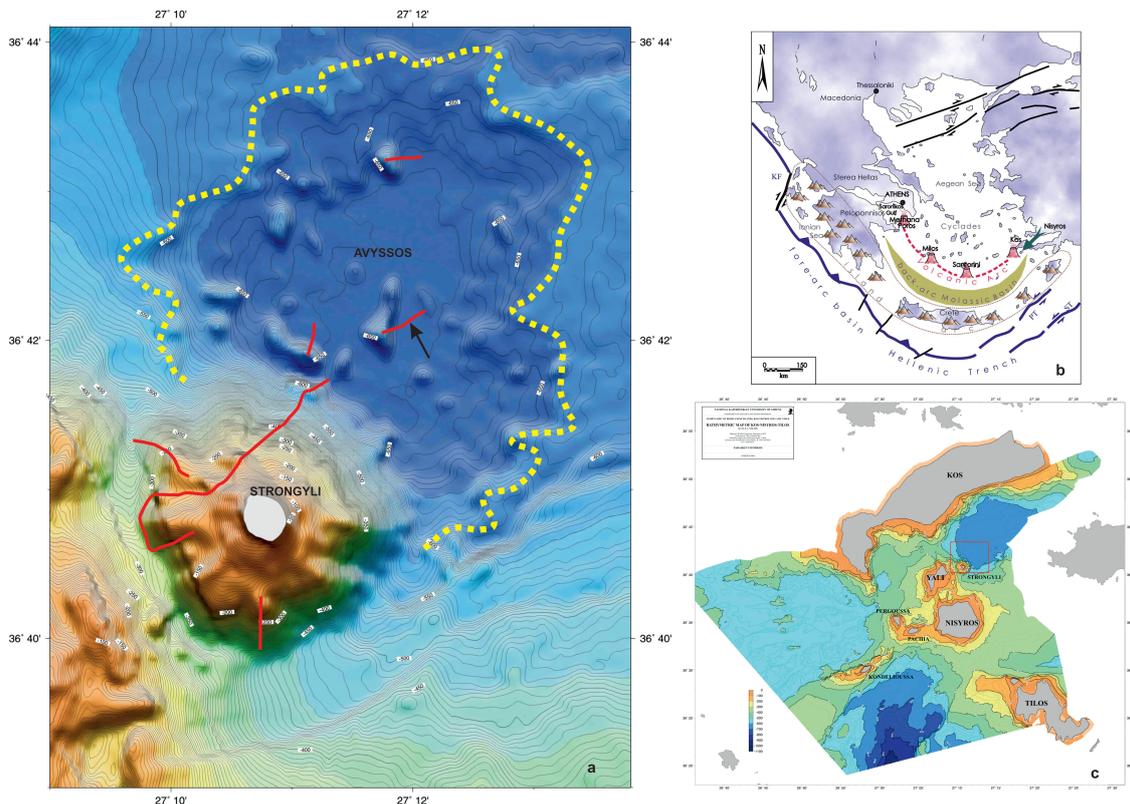


Figure 1. (a) A bathymetric map of Avysson off the coast of Nisyros (adopted from [8]). The yellow dashed line roughly defines the borders of Avysson caldera and the red lines the ROV transects. A black arrow indicates the location of the ROV during data sampling. (b) The geodynamic structure of the Hellenic Arc. (c) The extended Kos-Nisyros-Tilos volcanic field in the eastern Aegean Sea is shown in the bottom right inset.

An efficient method to study those characteristics is to employ ROV-based sensors, able to record CTD data (Conductivity, Temperature, Depth) in the water column. In the case of Kolumbo, the characteristics of the depth profiles of temperature, conductivity and salinity over the vent field were directly associated with the morphology of its crater [12]. In addition, Kolumbo's hydrothermal vent activity was recorded as time series of temperature and conductivity during two different periods; in 2010, the volcano exhibited intense activity, while in 2011, its activity decreased significantly. Besides the profound changes in the CTD data, a novel mathematical technique was applied for the first time on the time series to identify the underlying stochastic processes [14] that drive the activity of the volcano.

Due to the different characteristics and volcanic activity in comparison with Kolumbo, Avysson presents an ideal test case to examine the applicability of the mathematical method. The absence of hydrothermal vent field activity and the open caldera can serve as the exact opposite benchmarks for the mathematical technique as applied in Kolumbo, which can help identify underlying similarities and differences between the two submarine volcanic structures. As such, based solely on statistical

markers, the combined results from the two different submarine volcanoes can potentially serve as geohazard assessors for similar environments where no information is available.

2. Methodology

The CTD data were acquired by a pre-calibrated SeaBird FastCAT 49Plus CTD probe mounted on ROV “Hercules”, which was tethered to ROV “Argus” receiving power from the ship and transferring data via a fiber optics cable. The CTD sensor accumulated data continuously during dives at a 16-Hz rate [15]. No other sensors were connected to the CTD at the time of data recording. Detailed information on data collection by means of the CTD sensor aboard the ROV can be found in [11,12,14]. All data used in the statistical analysis described below were taken at the deepest location possible inside Avyssonos caldera with the ROV hovering at 1 m above the sea bottom at a fixed location, while the ship remained stationary at the surface. The time series of temperature, T , and conductivity, E_C , recorded under these conditions have been analyzed by assuming that they are manifestations of stochastic processes (see Figure 2).

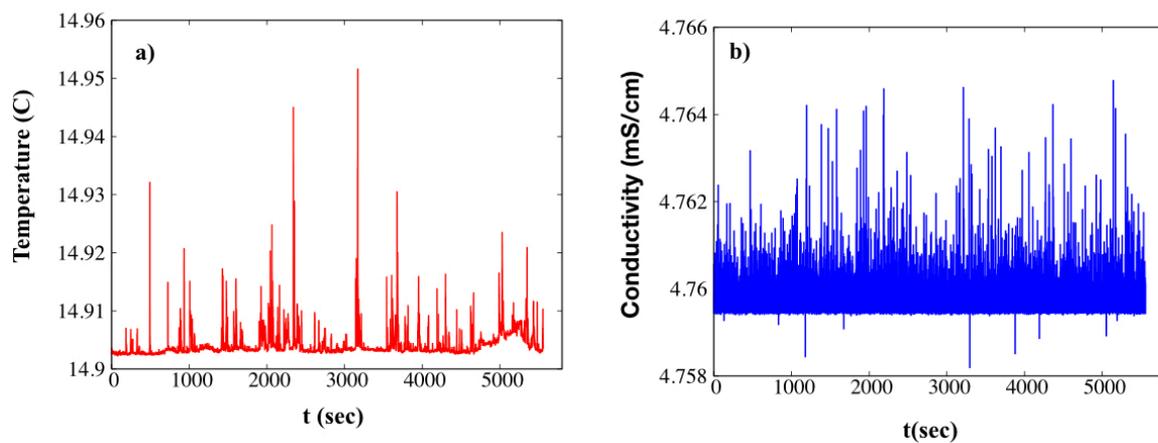


Figure 2. The time series (summer 2010) of (a) temperature and (b) conductivity. Notice that conductivity changes are very small and present a regularity; instead, temperature variations are a bit larger and possess some irregular peaks.

A first estimate of the behavior of a stochastic process, $x(t)$, can be made in terms of its variance $W(t) = \langle x^2(t) \rangle - \langle x(t) \rangle^2$. In general, the variance in time of a stochastic process scales as:

$$W(t) = \frac{K_\gamma t^\gamma}{\Gamma(1 + \gamma)} \tag{1}$$

where K_γ is a generalized coefficient expressed in proper units, e.g., if Equation (1) describes the variance of temperature, then the units are $\text{deg}^2 \cdot \text{s}^{-\gamma}$. Equation (1) can serve as a tool for a first classification of the stochastic process based on the value of the exponent γ . A stochastic process is classified as sub-normal or anti-persistent for $0 < \gamma < 1$, Brownian under certain conditions for $\gamma = 1$ (see Discussion below), super normal or persistent for $1 < \gamma < 2$, ballistic for $\gamma = 2$, and stationary for $\gamma = 0$.

For discrete data time series Equation (1) takes the form:

$$W(\Delta, T) = \frac{1}{M} \sum_{i=1}^M \frac{1}{T - \Delta} \sum_{n=1}^{T-\Delta} (x_i(n + \Delta) - x_i(n))^2 - \left(\frac{1}{M} \sum_{i=1}^M \frac{1}{T - \Delta} \sum_{n=1}^{T-\Delta} (x_i(n + \Delta) - x_i(n)) \right)^2 \tag{2}$$

where M is the number of experiments repeated under exactly the same conditions, so that Equation (2) expresses the ensemble averaged over all records, while for $M = 1$, the first summation is omitted.

T is the length of the recorded trajectory, and Δ is the lag time, which is the elapsed time between two measurements and plays the role of time for discrete datasets.

More insights on the mechanisms governing a stochastic process can be extracted by using sophisticated methods appropriate for time series analysis. Among them [16–21] the Generalized Moments Method (GMM) is generally the more robust and works well even for short time series [22]. It has been successfully applied in numerous fields [14,23–26]. Considering a time series in the form $x(n)$ with $n = 1, 2, \dots, N$, where N is the total number of steps (measurements), and if the minimum lag time defined as the reciprocal of the sampling rate, f , or in other words the time elapsed between two consecutive measurements is τ , then the total length of the trajectory is $T = N \times \tau$. The method works as follows:

- Step 1: Construction of time series with different lag times.

We create new time series, $y_n(\Delta)$, which contain the absolute change of the values between two points of the initial series that are apart by Δ .

$$y_n(\Delta) = (x(n + \Delta) - x(n))^2 \text{ for } n = 1, 2, \dots, T - \Delta \text{ and for } \Delta = \tau, 2\tau, 3\tau, \dots, \tau_{max}\tau.$$

In order to have statistically-reliable results, we define the maximum lag time as one tenth of the maximum length of the original time series, $\tau_{max} = \frac{N}{10}$, creating thus $N/10$ new time series of length $N - \Delta$ each.

- Step 2: Estimate of the statistical moments.

We estimate the moments of $y_n(\Delta)$ according to Equation (3):

$$\rho(q, \Delta) = \frac{1}{T - \Delta} \sum_{n=1}^{T-\Delta} (y_n(\Delta))^q \tag{3}$$

where even fractional values of the moment, q , are taken into account. We use only positive values of the moments [27]. The values of the moments in the range $0 < q \leq 2$ are responsible for the core of the probability density function (pdf), while moments higher than 2, $q > 2$, contribute to the tails of the pdf.

- Step 3: In general, the moments will scale according to Equation (4).

$$\rho(q, \Delta) \approx \Delta^{z(q)} \tag{4}$$

where $z(q)$ is the structure function whose shape gives information on the stochastic mechanism(s) governing the motion.

Stochastic processes can be of a varying degree of complexity. When the structure function takes the simple form $z(q) = Hq$, there is a direct relation of exponent γ of Equation (1) and the cofactor of the structure function, namely $H = 2\gamma$, where H is the Hurst exponent. If $z(q)$ is a linear function of q or is linear in different portions (bilinearity), then the mechanisms that likely govern the stochastic sequence are additively decomposable, e.g., fractional Brownian motion (fBm), fractional Gaussian noise (fGn), Lévy flights and/or walks. Linear dependency of $z(q)$ on q describes monofractal processes. Instead, if $z(q)$ has a convex shape as a function of the order of the moment, q , then multiplicative processes drive the stochastic sequence, which also means that H is a function of q . For some values of q , the exponents $H(q) = z(q)/q$ are associated with special features. For $q = 1$, the value of $H(1)$ describes the scaling behavior of the absolute value of the increments and is identified with the Hurst exponent if the process is monofractal. For $q = 2$, the value of $H(2)$ is associated to the scaling exponent of the autocorrelation function and is related to the power spectrum [28]. Of all types of multifractals, universal multifractals are likely to be ubiquitous [29–31], and the structure function reads [31]:

$$z(q) = Hq - \frac{C}{a-1}(q^a - q) \tag{5}$$

For $a = 1$, Equation (5) takes the form $z(q) = Hq - Cq \log(q)$, and the distribution draws changes from a Cauchy–Lorentz distribution. For $a = 2$, Equation (5) reads $z(q) = Hq - C(q^2 - q)$, and the distribution draws changes from a lognormal distribution.

3. Results

We estimated the variance for both the temperature and conductivity time series depicted in Figure 2, and the results are illustrated in Figure 3. The variance for the conductivity time series (right panel in Figure 3) was constant as a function of the lag time, Δ . This behavior indicates that conductivity measurements describe a stationary process, which in the present case by also combining results from moments analysis (see below) pointed to a submarine volcano with no fingerprints of any hydrothermal vent field activity. In contrast, the temperature time series showed a power-law behavior for small lag times followed by an almost stationary regime for longer lag times.

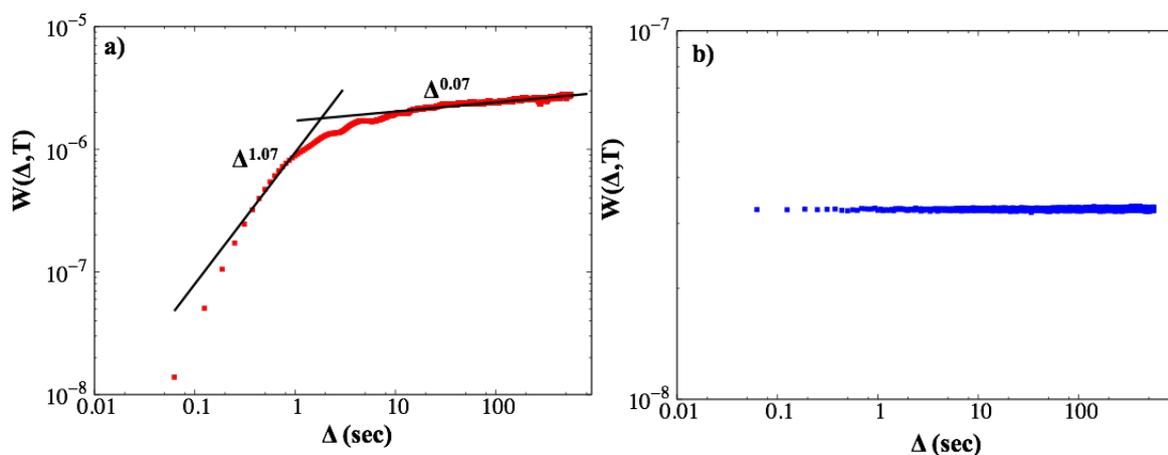


Figure 3. The logarithmic plots of variance estimated by using Equation (2) versus the lag time for (a) temperature and (b) conductivity.

The variance of temperature is formed by two discrete regimes (left panel in Figure 3), with a turning point at $\Delta = 1$. We fit the variance of temperature with a function of the form, $W(\Delta) = b\Delta^\gamma$. For the first regime, $1/16 < \Delta < 1$, we found that $\gamma = 1.073 \pm 0.051$, and for the second regime, the value of the exponent was $\gamma = 0.0750 \pm 0.0003$. In the first regime, the obtained value of 1.073 indicates a process that potentially can be characterized as Brownian, since $\gamma \approx 1$. However, the form of the structure function (Figure 4) does not support this hypothesis (see also the discussion in [26]) and reveals information useful for characterizing the system, as closed or open, i.e., a bowl-shaped caldera vs. a caldera with no large difference between its rim and its bottom, respectively. In the second regime, the estimated exponent showed clearly a stationary process, suggesting that all measurements which were at least 1 s apart shared the same properties.

We estimated the moments for both temperature and conductivity by using Equation (3). The results are illustrated in Figure 5. For conductivity, all moments were constant, confirming the stationarity of the recorded data and indicating that the volcano was at rest (Figure 5b). By comparing this to Kolumbo's rest period, where the conductivity time series are multifractal, a result of the outgoing ions and gases from the hydrothermal field, one can conclude that a hydrothermal field is rather absent in this case. For temperature, all the moments showed the existence of two distinct regimes with a turning point at $\Delta = 1$ s. For lag times greater than 1 s, all moments were parallel to the time axis, meaning that the temperature recorded at a fixed location over the volcanic field had been equilibrated with the temperature of the surrounding water. For lag times smaller than 1 s, the moments of temperature have been fitted by Equation (4). The obtained exponents were used to find the form of the structure function (see Figure 4). Every point (red squares) depicted in Figure 4 is

the value of the structure function for each moment. The structure function, $z(q)$, posed a convex shape as a function of the order of the moment, q , suggesting the existence of a multifractal process. In other words, multifractality underlined the existence of two or more stochastic processes acting in a multiplicative way and defining the overall behavior of the temperature time series for small lag times.

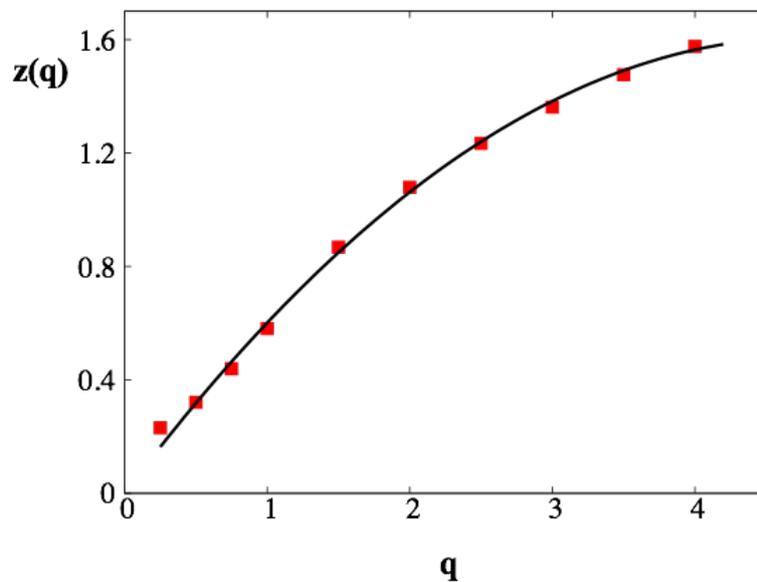


Figure 4. The structure function $z(q)$ versus the order of the moment q for the temperature time series.

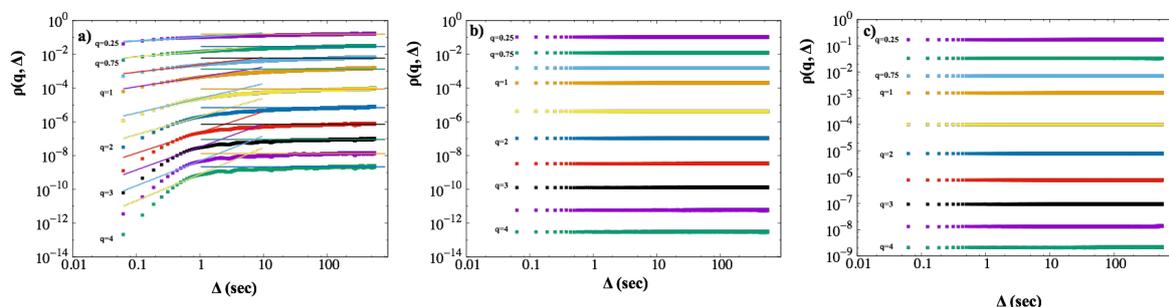


Figure 5. Generalized moments of (a) temperature and (b) conductivity as functions of the lag time, Δ , for different moments q ; (c) the same as (a) for the temperature shuffled time series.

The data presented in Figure 4 are best fitted by Equation (4) for $a = 2$, lognormal distribution. The parameters H and C had values of $H = 0.601 \pm 0.011$ and $C = 0.070 \pm 0.005$. A lognormal distribution implies the existence of a multiplicative effect, which is the result of at least two random processes. In our case, such an effect may have been caused by a random process causing temperature variations inside the crater and a random process related to the variations of the temperature of the surroundings. The overall effect says that the volcano area operates as an open system where a rapid mixing (first regime) between the outflushing hydrothermal fluid from the vents (if any) with its surroundings washes out any small temperature variations and drives the volcanic field to equilibrium (stationary second regime). In the first regime, the observed multifractality was likely due to (a) a broad distribution of the probability values and (b) long-range correlations. Randomization of the original time series to produce a shuffled time series can give insights into the origin of multifractality. If GMM analysis retains multifractal behavior for the shuffled time series, then the broad distribution of probability values is the origin of multiple and non-decomposable time scales. In this case, the system

tries to maximize information. Instead, if for the shuffled time series, the GMM analysis shows random or stationary behavior, existing long-range correlations are responsible for multifractality. Performing the GMM analysis for the shuffled temperature time series, we found that all moments are constant and independent of the lag time; see Figure 5c. Long-range correlations are responsible for the observed dynamics of temperature in the first regime. These correlations were likely caused by the fast mixing of the two different fluid-banks.

Analyzing the recorded time series of temperature and conductivity just above the Avyssos submarine volcano, we found that (a) the volcano was at rest and there was no activity in terms of emerging gases and ionic particles and (b) the temperature presented an initial multiplicative behavior reflecting the mixing of at least two different random processes, which then turns to be stationarity.

The early-time temperature behavior is an indicator that the Avyssos submarine volcano acts as an open system, and this observation is in line with geomorphological mapping of the volcano [8] and visual observations with the ROV that showed no activity [15]. In the case of the Kolumbo vent field, where the same methodology has recently been applied to analyze temperature and conductivity recorded over the active vents [14], it was observed that when Kolumbo was at rest, the conductivity for early times presented a multifractal behavior mirrored by the outcoming gases and ions from the vent field, while the temperature was a stationary process for all times, thus indicating that the system operated as a closed one. The significance in the context of the present work is the fact that the combined results provide the two extremes of the ranges showing that the GMM can be applied to characterize any other system, open or closed. Based on these findings, we propose the above methodology of analysis as an indicator for the state of activity and the shape (closed or open systems) of a submarine volcano featuring a volcanic field. On the one hand, when conductivity presents a multifractal character, it also serves as an indication of vent activity, while if conductivity behaves as a stationary process, then activity is not present. On the other hand, if temperature displays a stationary behavior, independent of the lag time, the system operates as a closed one in relation to its surroundings, while if multifractality is observed for small lag times, the system is an open one.

4. Conclusions

The temperature and conductivity time series recorded by a ROV just above the seafloor at the underwater caldera of Avyssos have been analyzed with the GMM method. The analysis showed that when the volcano is at rest, conductivity is a stationary process, and the crater of the volcano is in direct contact with the overlying water environment, thus defining an open system, as shown by the temperature early lag time behavior. Comparing the results of the same method with the case of the Kolumbo submarine volcano, where conductivity showed the existence of multiple time scales and temperature was stationary, we suggest that the behavior of these two properties can serve as an indicator of the state of activity (conductivity) and the physical setting, open or closed, of the submarine volcano (temperature). The present methodology can be useful especially when insights are needed for poorly-known deep-sea volcanoes, where in situ measurements by means of the use of submarine robotic technologies are either too expensive or even technically impossible. Further application of the technique to other volcanic systems is required to confirm its applicability as a universal geohazard warning method.

Author Contributions: E.B. and T.J.M. contributed to the data preparation and mathematical analysis. P.N. led the process of data sampling and map reconstruction. E.B., T.J.M., P.N. and F.Z. contributed to the manuscript preparation, data assessment and interpretation.

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Conflicts of Interest: The authors declare no conflict of interest.

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